

# Calculation of Lamb shift for states with $j = 1/2$

V. Patkóš<sup>1</sup>, J. Zamastil<sup>1</sup>, and D. Šimsa<sup>1</sup>

<sup>1</sup> *Department of Chemical Physics and Optics, Faculty of Mathematics and Physics, Charles University, Prague*

Presenting Author: simsadan@gmail.com

The Lamb shift in the light hydrogen-like atom yields one of the most stringent test of quantum electrodynamics. It is customary to write the Lamb shift of general  $S$ -state in the form

$$\Delta E_n = \frac{n^3 \Delta E_n - \Delta E_1}{n^3} + \frac{\Delta E_1}{n^3}, \quad (1)$$

where  $n$  is the principal quantum number of the state under consideration. The first and the second terms on the right-hand side are referred to as the state-dependent and the state-independent parts, respectively. The state-dependent part of the  $S$ -states is nearly completely given by the self-energy effect in one-loop approximation which is usually expressed in the form

$$\Delta E_n = \frac{\alpha (Z\alpha)^4}{\pi n^3} F(n, l_j, Z\alpha), \quad (2)$$

where  $l$  and  $j$  are the orbital and total angular quantum numbers of the state under consideration,  $Z$  is the charge of the nucleus in units of elementary charge  $e$ ,  $\alpha = e^2/(4\pi)$  is the fine structure constant.

It was suggested in [1,2] to calculate the dimensionless function  $F$  by means of relativistic generalization of multipole expansion (RME)

$$F = \sum_{v=1}^{\infty} F_v, \quad (3)$$

where  $F_v$  are the "relativistic multipoles" (for formulas see Eq. (73) of [2]).

Here, we report the evaluation of the state-dependent part of the  $S$ -states and  $P_{1/2}$  states by means of RME. We obtain the results for  $Z = 1 - 50$  and  $n = 2 - 10$ .

The results for low  $Z$  are the most accurate results given so far in the literature. In the case of hydrogen the uncertainty is significantly less than 1 Hz. Some of the results are presented in the following table.

| Term  | State     | $Z = 1$                   | $Z = 2$                   | $Z = 3$                   | $Z = 4$                   |
|-------|-----------|---------------------------|---------------------------|---------------------------|---------------------------|
| Lead  | $2s - 1s$ | 0.229991606931            | 0.230390709933            | 0.230936267290            | 0.231594004269            |
|       | $F_3$     | 0.000039870858            | 0.000154033610            | 0.000336061748            | 0.000581009986            |
|       | $F_4$     | $6.241492 \times 10^{-8}$ | $4.634566 \times 10^{-7}$ | $1.469760 \times 10^{-6}$ | $3.298274 \times 10^{-6}$ |
| Total |           | 0.2300315402(3)           | 0.230545207(2)            | 0.231273799(6)            | 0.23217831(1)             |
| Lead  | $3s - 1s$ | 0.288771400893            | 0.289100037681            | 0.289541947085            | 0.290067435815            |
|       | $F_3$     | 0.000048953133            | 0.000189443771            | 0.000413949977            | 0.000716666402            |
|       | $F_4$     | $7.723537 \times 10^{-8}$ | $5.752828 \times 10^{-7}$ | $1.829313 \times 10^{-6}$ | $4.115199 \times 10^{-6}$ |
| Total |           | 0.2888204313(3)           | 0.289290057(2)            | 0.289957726(8)            | 0.29078822(2)             |
| Lead  | $4s - 1s$ | 0.312542251722            | 0.312795865525            | 0.313129873756            | 0.313519935279            |
|       | $F_3$     | 0.000052358450            | 0.000202772029            | 0.000443363221            | 0.000768038994            |
|       | $F_4$     | $8.348329 \times 10^{-8}$ | $6.226729 \times 10^{-7}$ | $1.982136 \times 10^{-6}$ | $4.463102 \times 10^{-6}$ |
| Total |           | 0.3125946937(3)           | 0.312999260(3)            | 0.313575219(8)            | 0.31429244(2)             |

Contribution of individual multipoles to the normalized difference of  $S$ -states for  $n = 2 - 4$ . Lead stands for first two multipoles,  $F_1 + F_2$ .

## References

- [1] J. Zamastil, V. Patkóš, Phys. Rev. A **86**, 042514 (2012).  
 [2] J. Zamastil, V. Patkóš, Phys. Rev. A **88**, 032501 (2013).